

# Numerical integration methods (ODE+PDE)

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# Numerical integration of differential equations

- We search a solution (a numerical approximation) for an ordinary differential equation

$$\frac{d^2 y}{dt^2} = -ky$$

- Or a partial differential equation:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$$

- Given appropriate initial and boundary conditions

# Numerical integration of ODEs

- NB: any ODE of order  $> 1$  can be written as a system of first order equations Eg:

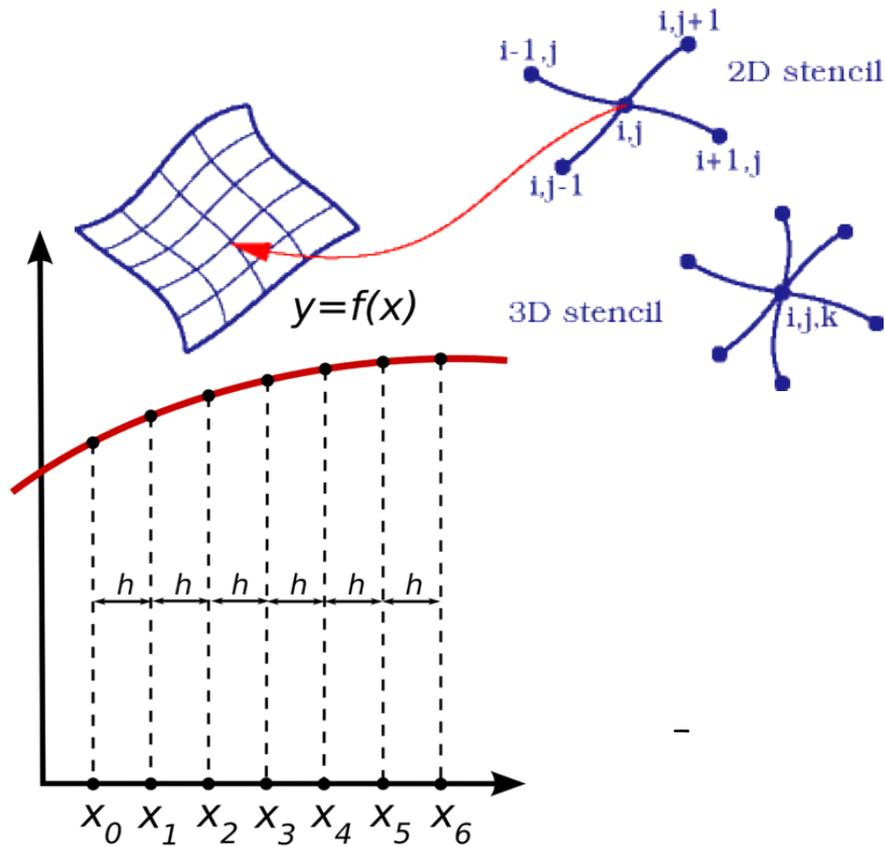
$$\frac{d^2 y}{dt^2} = -ky \quad \longrightarrow \quad \begin{cases} \frac{du}{dt} = -ky \\ \frac{dy}{dt} = u \end{cases}$$

- General problem:

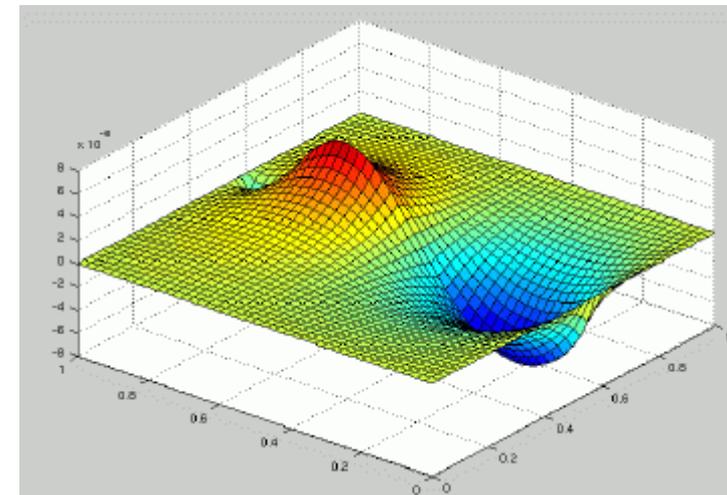
$$\frac{dy_i}{dt} = f_i(y_1, y_2 \dots y_n; t)$$

# Finite difference methods

- We substitute to the continuous problem a representation on a discretized grid (in space and time)

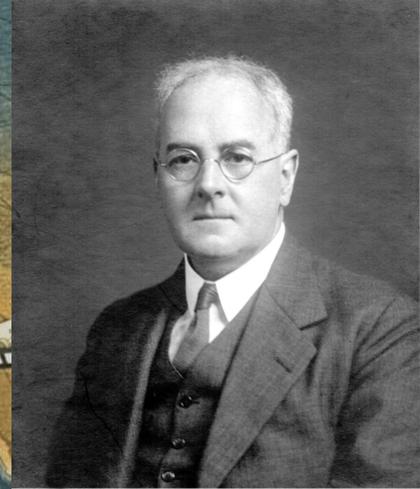


$$f'(a) \approx \frac{f(a+h) - f(a)}{h}$$

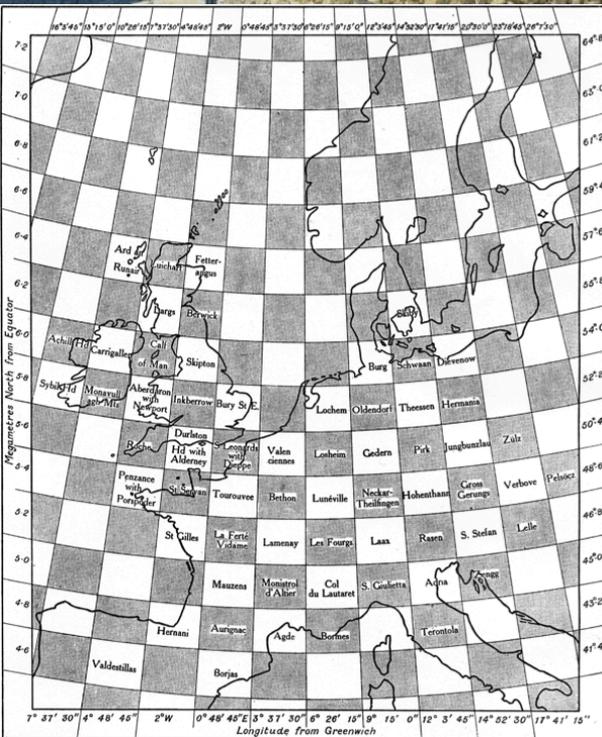


# L.F. Richardson, Weather Prediction by Numerical Process (1922):

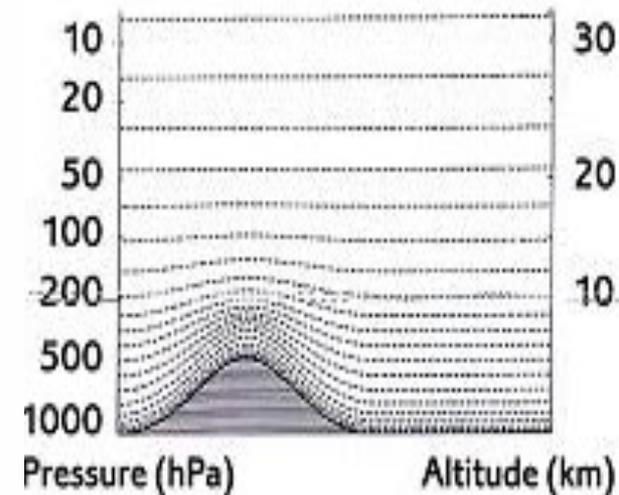
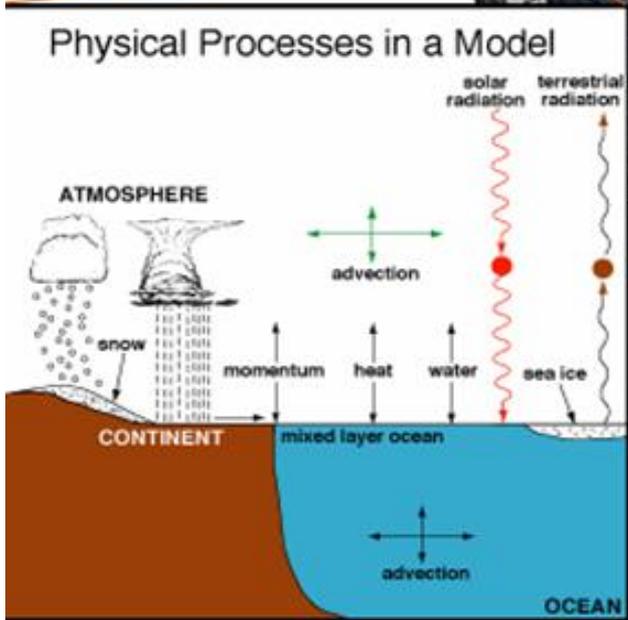
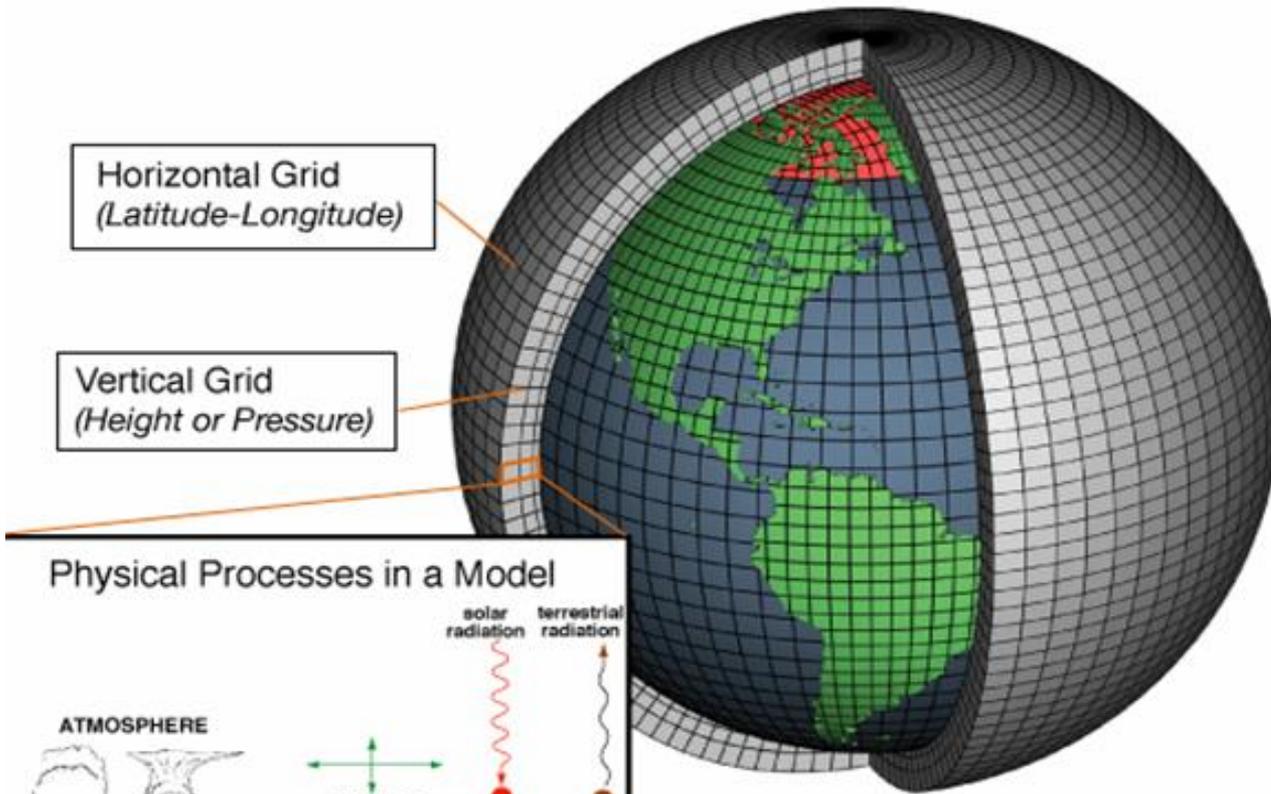
## The forecast factory



L. F. Richardson, 1931



# Numerical circulation models



Mathematical equations that represent the physical characteristics and processes are entered for each box

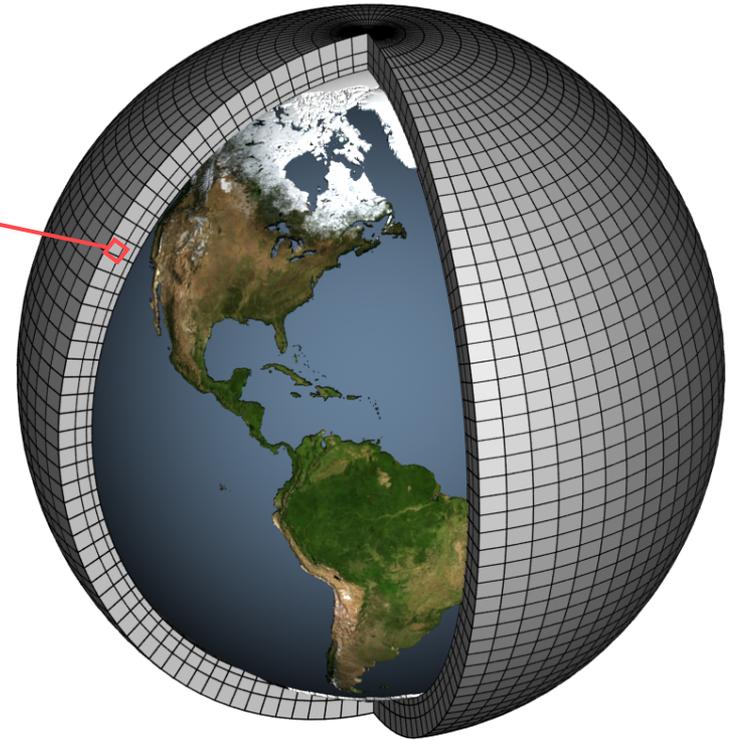
$$\frac{\partial u}{\partial t} = \eta v - \frac{\partial \Phi}{\partial x} - c_p \theta \frac{\partial \pi}{\partial x} - z \frac{\partial u}{\partial \sigma} - \frac{\partial \left( \frac{u^2 + v^2}{2} \right)}{\partial x}$$

$$\frac{\partial v}{\partial t} = -\eta \frac{u}{v} - \frac{\partial \Phi}{\partial y} - c_p \theta \frac{\partial \pi}{\partial y} - z \frac{\partial v}{\partial \sigma} - \frac{\partial \left( \frac{u^2 + v^2}{2} \right)}{\partial y}$$

$$\frac{\delta T}{\delta t} = \frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z}$$

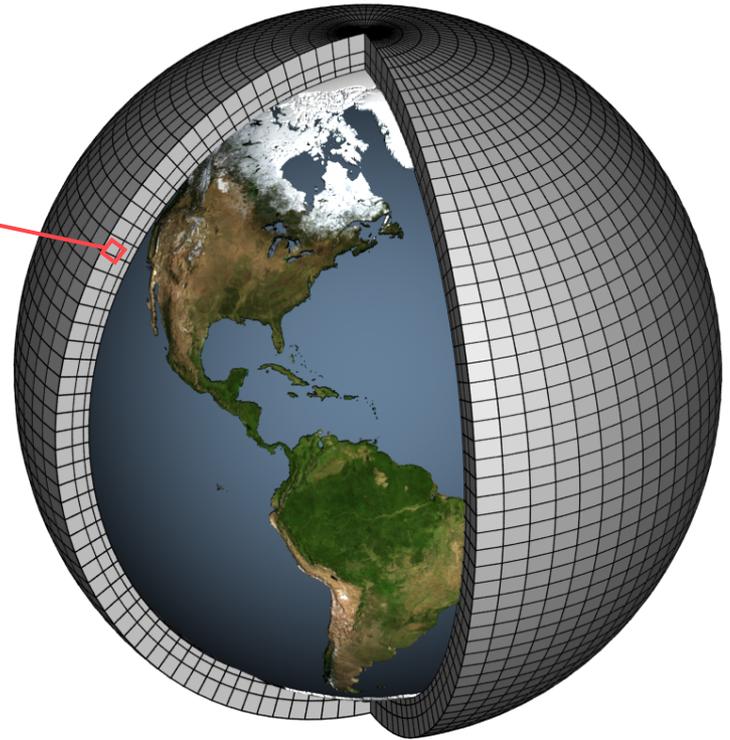
$$\frac{\delta W}{\delta t} = u \frac{\partial W}{\partial x} + v \frac{\partial W}{\partial y} + w \frac{\partial W}{\partial z}$$

$$\frac{\partial}{\partial t} \frac{\partial p}{\partial \sigma} = u \frac{\partial}{\partial x} x \frac{\partial p}{\partial \sigma} + v \frac{\partial}{\partial y} y \frac{\partial p}{\partial \sigma} + w \frac{\partial}{\partial z} z \frac{\partial p}{\partial \sigma}$$



# Equations are converted to computer code and climate variables are set

```
if (diagts .and. eots) then
  do 1500 m=1,nt
    do 1490 k=1,km
      fx = cst(j)*dym(j)*dzt(k)/(c2dtts*dtxcel(k))
      do 1480 i=2,imtml
        boxfx          = fx*dxt(i)*fm(i,k,jc)
        sddt           = (ta(i,k,m)-t(i,k,jc,nm,m))*boxfx
        svar           = (ta(i,k,m)**2-t(i,k,jc,nm,m)**2)
                      *boxfx
        n              = 0
        termbt(k,1,m,n) = termbt(k,1,m,n) + sddt
        tvar(k,m,n)    = tvar(k,m,n)      + svar
        n              = nhreg*(mskvr(k)-1) + mskhr(i,j)
        if (n .gt. 0 .and. mskhr(i,j) .gt. 0) then
          termbt(k,1,m,n) = termbt(k,1,m,n) + sddt
          tvar(k,m,n)    = tvar(k,m,n)      + svar
```



# Finite differences representation of derivatives

Can be obtained from

- Classical definition of first order derivative of a function  $u(x,y)$  at a point:

$$\frac{\partial u}{\partial x} \Big|_o = \lim_{\Delta x \rightarrow 0} \frac{u(x_o + \Delta x, y_o) - u(x_o, y_o)}{\Delta x}$$

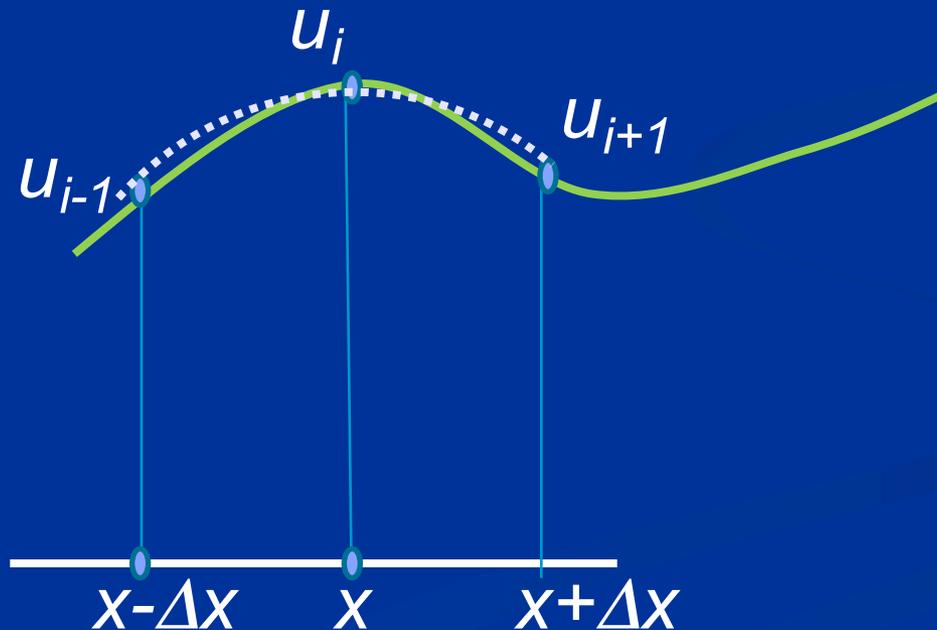
- Taylor expansion of  $u(x,y)$  around a point

$$u(x_o + \Delta x, y_o) = u(x_o, y_o) + \Delta x \frac{\partial u}{\partial x} \Big|_o + \frac{\Delta x^2}{2} \frac{\partial^2 u}{\partial x^2} \Big|_o + O(\Delta x^3) \dots$$

# Finite differences representation of derivatives

Can be obtained from:

- Fitting a polynomial around a point::



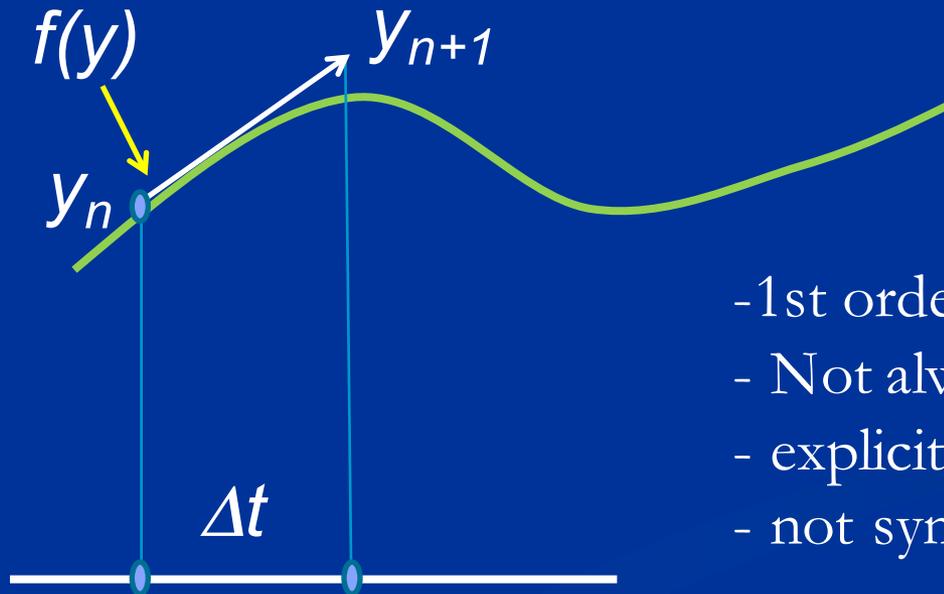
# ODE: The Euler method

$$\frac{dy}{dt} = f(y)$$

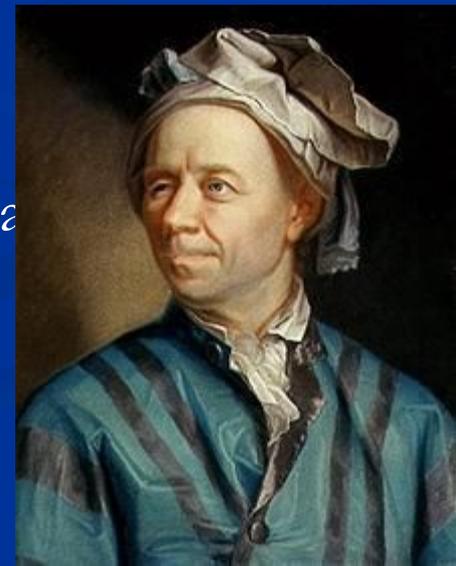
$$y(t + \Delta t) = y(t) + \Delta t f(y(t))$$

$$y_{n+1} = y_n + \Delta t f(y_n)$$

+ Err.ore  
Troncamento



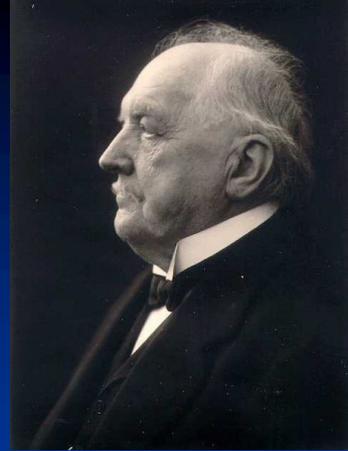
- 1st order method accurate
- Not always stable
- explicit
- not symmetric





C.D.T. Runge

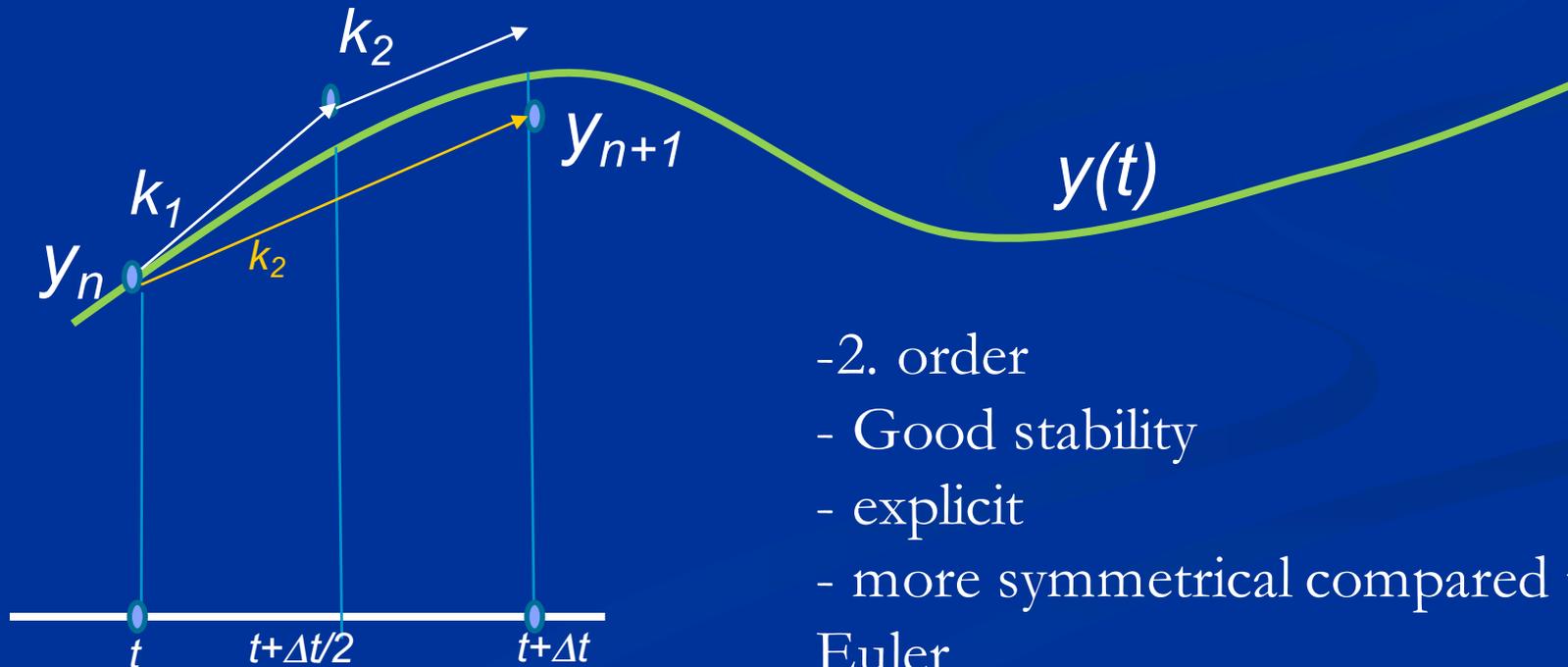
# ODE: Runge-Kutta 2



M. W. Kutta

$$\frac{dy}{dt} = f(y)$$


$$\left\{ \begin{array}{l} k_1 = f(y) \\ k_2 = f\left(y + \frac{\Delta t}{2}k_1\right) \\ y(t + \Delta t) = y(t) + k_2\Delta t + O(\Delta t^3) \end{array} \right.$$



- 2. order
- Good stability
- explicit
- more symmetrical compared to Euler

# ODE: Runge-Kutta 4

$$\frac{dy}{dt} = f(y)$$

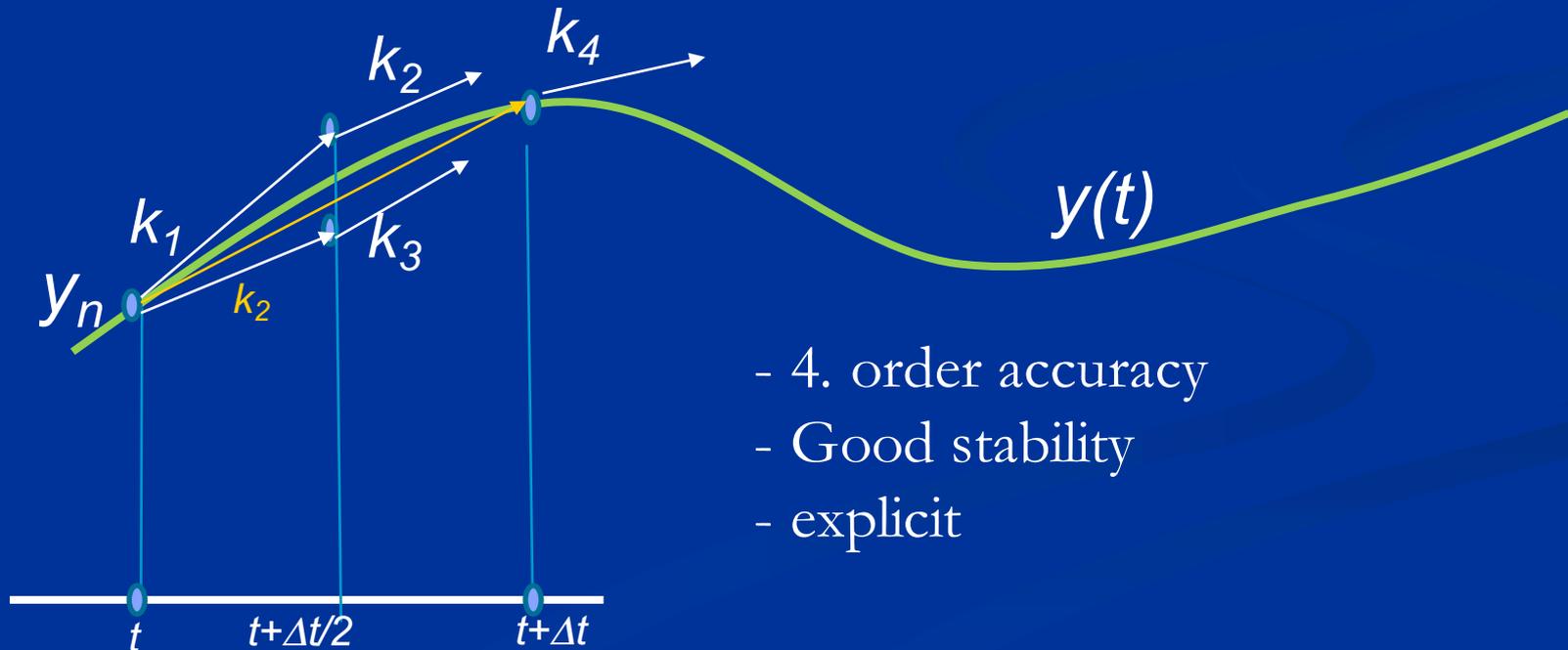
$$k_1 = f(y)$$

$$k_2 = f\left(y + \frac{\Delta t}{2}k_1\right)$$

$$k_3 = f\left(y + \frac{\Delta t}{2}k_2\right)$$

$$k_4 = f\left(y + \Delta tk_3\right)$$

$$y(t + \Delta t) = y(t) + \left[\frac{k_1}{6} + \frac{k_2}{3} + \frac{k_3}{3} + \frac{k_4}{6}\right] \Delta t + O(\Delta t^5)$$



- 4. order accuracy
- Good stability
- explicit

# Other methods

- Leapfrog
- .... Many other explicit methods
- Predictor-corrector
- Implicit methods (greater stability, non necessarily more accurate)

# PDE, examples

$$\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial x} = 0$$

Scalar advection

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = -v \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x}$$

FT

CS

Non è stabile!

$$\frac{\partial u}{\partial t} = \nu \frac{\partial^2 u}{\partial x^2} + \nu \frac{\partial^2 u}{\partial y^2}$$

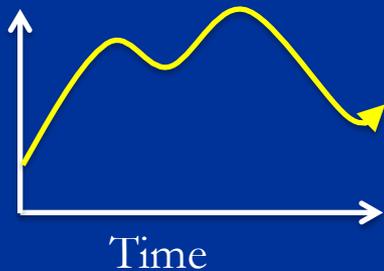
Heat equation

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = \nu \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{\Delta x^2}$$

# Initial condition problems vs. boundary problems

## ■ Initial conditions problems:

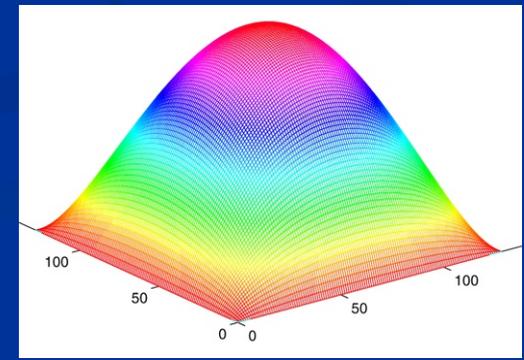
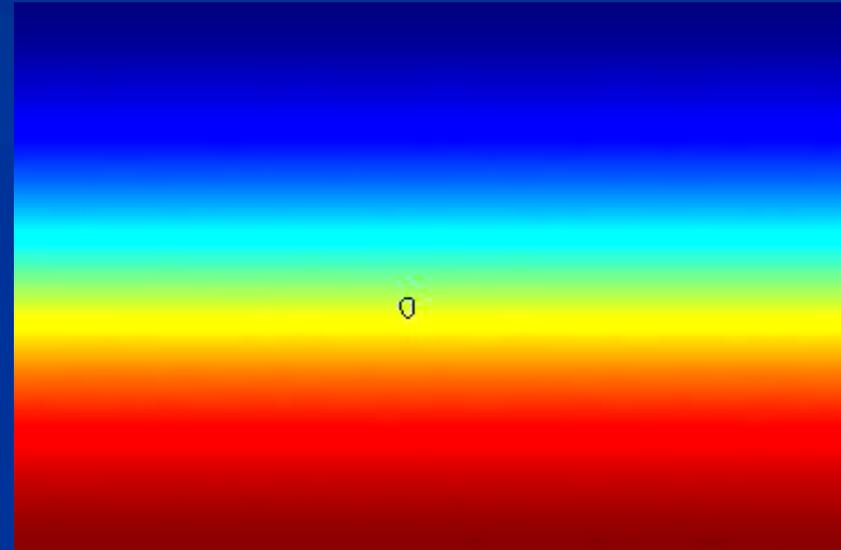
- Choice of initial conditions
- Evolution equations
- Boundary conditions



## ■ Boundary problems

- Equations to be solved in the domain
- Boundary conditions

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \rho$$



# PDE: Boundary conditions

- Dirichlet conditions.  $u=f$  su  $\partial\Omega$
- von Neumann conditions:  
eg:  $\partial u/\partial n=f$  or  $\partial u/\partial s=g$  in  $\partial\Omega$
- Mixed conditions e.g:  $\partial u/\partial n+ku=f$



